SHORT COMMUNICATION

Intense Upconversion Luminescence of Yb³⁺-Er³⁺ in Li₂O Content Tungsten-tellurite Glasses

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Abstract The Er³⁺ -Yb³⁺ codoped in Li₂O content tungsten -tellurite (TWL) transparent glasses are synthesized and measured the absorption, Raman and upconversion luminescence (UPL) spectra. At room temperature intense green emission peak at 560 nm (${}^{4}S_{3/2} \rightarrow {}^{4}I_{15/2}$) and red emission peak at 670 nm (${}^{4}F_{9/2} \rightarrow {}^{4}I_{15/2}$) of Er³⁺ observed even at minimum 86 mW pumping power of infrared 980 nm excitation. For structure of the TWL glass, Raman spectrum result revealed that an important role of WO₃ in the formation of glass network linkage with Li₂O. Under this influence estimated lifetime of the ${}^{4}I_{11/2}$ of Er^{3+} was 1.89 µs and due to lower phonon energy of the glass produce strong upconversion signal. The effect of Er₂O₃ concentration on emission intensity result indicated that green emission intensity initially increase in compare to red emission. Under the 980 nm pump power variation measured the relatively increases the red emission to the green emission intensity and analyze the possible upconversion mechanism and process.

Keywords Lanthanide ion-doped tellurite glasses · Optical properties erbium · Upconversion luminescence · Raman spectrum

Introduction

The trivalent rare-earth ions such as Er³⁺, Tm³⁺,Ho³⁺, Pr³⁺ and Nd³⁺ doped materials have scientific and techno-

G. F. Ansari · S. K. Mahajan (⊠) · J. Parashar Department of Applied Physics, Samrat Ashok Technological Institute, Vidisha, MP 464001, India e-mail: sachin k mahajan@rediffmail.com logical interest due to their optical properties in developing photonic devices, compact laser and broad band amplifier [1-3]. Among the rare-earth ions, Er^{3+} is most popular as well as most efficient ions. In addition with Yb³⁺ ions as a sensitizer, perform an important role to improve the energy transfer to Er³⁺ ions and luminescence efficiency accordingly [4, 5]. The host material for rare-earth ions play an important role in obtaining efficient upconversion signal, since glass host with low phonon energy can reduce nonradiative loss to multiphonon relaxation and thus realize strong upconversion luminescence [6-8]. Generally, upconversion is difficult to observe in observe in oxide crystals or glasses compared with that fluoride glasses. But oxide glasses might be much better for practical applications because of their high thermal stability, chemical durability and ease of fabrication [9, 10].

In recent years, tellurite glasses are growing interest due to their low phonon energy of 600-800 cm⁻¹, which is significantly lower than that of both silicates $(1,100 \text{ cm}^{-1})$ and germinates (880 cm⁻¹), and the large refractive index (>2.0) enhances both the absorption and emission cross sections [11, 12]. Also tellurite glasses are more chemically and environmentally stable than fluoride glasses, and have high rare-earth ions solubility. This has promoted a widespread interest in the spectroscopy investigation of tellurite glasses, especially the composition based on tungsten tellurite [13]. Comparing with sodium-tellurite and zinc-tellurite glasses have relatively low Tg and but better thermal stability which will helpful for fiber drawing [14]. Further crystals or glasses with addition of certain alkalie oxide as the network modifier could enhance the optical properties. Chen et al. observed enhanced upconversion emission in Y₂O₃ nanocrystal by codoping with Li^{3+} ions [15]. Shen et al. reported that Er^{3+} in tungstentellurite glasses mainly in the TeO₂-WO₃-R₂O(R=K, Na,

Li) glasses for 1.5 μ m emission for optical amplifier [16]. No publications concerned with upconversion fluorescence of Er³⁺/Yb³⁺ in Li₂O-WO₃-TeO₂ glasses have come to our notice, though structural study in different composition for same content of this glasses by X-ray photoelectron spectroscopy and Raman spectra have been reported [17, 18].

This work deals with the synthesis, structural and optical measurements of glass samples of Er^{3+}/Yb^{3+} codoped 68TeO₂-15Li₂CO₃-13WO₃ (TWL) system. The 15 mol% content of Li₂O was chosen because Na₂O-TeO₂ glasses show marked stability which is compositional-dependent with maximum at 20 mol% fraction of Na₂O [19]. Green and red emission of glass samples under the 980 nm excitation was measured to verify visible upconversion via energy transfer process from Yb³⁺ to Er^{3+} and possible upconversion mechanisms are discussed.

Experimental

The basic composition of the glass sample was (70-x)TeO₂- $15Li_2CO_3$ -13 WO₃- 2Yb₂O₃ (in mol%) where x range mixed with Er₂O₃ (Fluka 99.99%) concentration from 0.5 to 2 mol%. For undoped lithium tungsten tellurite glass sample composition was 70TeO₂-15Li₂CO₃-15WO₃ prepared for the Raman spectrum. All the starting constituents materials were purer than 99.9 % (TeO₂ Aldrich, Li₂CO₃ Otto, WO₃ Otto and Yb₂O₃ Fluka) in batch of 10gm was mixed in agate mortar and melted in alumina crucible with lid at temperature range 850-950 °C about 30-45 min in muffle furnaces at air atmosphere. The homogeneous and transparent form of glass disc of 6 mm in diameter and 1-1.5 mm thickness were obtained by pouring quickly the glass liquid into preheated stainless steel sheet plates. As made glasses were annealed below the glass transition temperature at 310 °C for 8 h and allowed to cool at room temperature to obtain transparent light pink colour samples for subsequent measurements.

The density of TWL: Er/Yb glass sample was measured according to the Archimedes's principle using water as an immersion liquid, to be 4.87gm cm⁻³. The refractive index was measured by Brewster angles method using a spectrometer with sodium –vapor lamp is 2.1. The Raman spectrum of the undoped TWL glass sample was measured by Raman spectrophotometer (Andor Technology SR370I) with diode laser at 758 nm was used as an incident light source. The visible and near –infrared absorption spectra of the glass sample was measured at room temperature using atomic absorption spectrophotometer(GBC Australia) over spectral range of 300–850 nm. The upconversion emission spectra were obtained with SEPEX Fluorolog 1680 (0.22 m double beam) spectrometer upon 980nm diode

laser (Scitec Instrument, UK) excitation adjustable to maximum power of 1 W. In order to compare the upconversion luminescence intensity of different Er_2O_3 with fixed Yb₂O₃ doped TWL glass samples as accurate as we can the position, power (200 mW) of pumping beam and width of slit to collect the luminescence signal were fixed to same condition and the sample was set at the same place in the experimental setup.

Results and Discussions

Absorption Spectrum

Figure 1 shows the absorption spectra for TWL: Er/Yb glass sample at room temperature. The absorption spectra of Er^{3+} ion consists of eight absorption bands centered at 800 nm,650 nm, 540 nm, 520 nm, 490 nm,451 nm,406 nm, 377 nm corresponding to the absorption from the ground state ${}^{4}I_{15/2}$ to the excited states ${}^{4}I_{9/2}$, ${}^{4}F_{9/2}$, ${}^{4}S_{3/2}$, ${}^{2}H_{11/2}$, ${}^{4}F_{5/2}$ + ${}^{4}F_{3/2}$, ${}^{2}\text{H}_{9/2}$, ${}^{4}\text{G}_{11/2}$ respectively. The absorption bands below 377 nm cannot be observed because of the band gap of the glass matrix and spectroscopic measurement was limited due to upper limit of spectrophotometer at 850 nm. However intense absorption band at the wavelength region from 950 nm to 1050 nm may be observed due to the large contribution of the absorption of Yb³⁺and Er³⁺, which arise from the transition ${}^{2}F_{7/2} \rightarrow {}^{2}F_{5/2}$ and ${}^{4}I_{11/2} \rightarrow {}^{4}I_{15/2}$ respectively as the band positions for Er³⁺-Yb³⁺ are similar to halide modified tellurite glasses [20]. Thus the samples present a good transparency in visible region indicating that possibility of practical applications in the visible region and negligibility of energy transfer from Er^{3+} to the glass matrix.



Fig. 1 Absorption spectrum of 68TeO₂-15Li₂CO₃-13WO₃: 2Er-2Yb (mol%) glass

Raman Spectrum of Undoped TWL Glass

Figure 2 shows the room temperature fluorescence subtracted Raman spectrum of undoped TWL glass. The Raman spectrum indicates the phonon frequency band peaked at 760 cm⁻¹ and 930 cm⁻¹ correspond to TeO₂ and WO₃ vibration respectively. The 760 cm^{-1} attributed to that of TeO₃ trigonal pyramid (tp) units due to stretching vibrations of TeO₃ and /or TeO₃₊₁. The highest peak at about 930 cm⁻¹ is observed corresponds to the stretching vibrations of W=O and W-O⁻ bonds associated with WO₄ and WO_6 groups. This mean that addition of WO_3 into TeO₂ result in the reduction of Te-O-Te linkages, and the formation of W-O-W and W-O-Te linkages [21]. These bands are similar and nearly same position of peaks was reported in the Raman spectrum of 85.5TeO₂-12.5WO₃ glass [22, 23]. But in our measurement side band near about 600 cm⁻¹ could not be observed due to lower limit of used Raman Spectrometer measurements. However, the Raman vibration at 460 cm⁻¹ is attributed to Te-O-Te chain unit symmetric stretching mode, while the spectral features from 610 cm⁻¹ to 680 cm⁻¹ and 760 cm⁻¹ correspond to TeO_4 bi-pyramidal arrangement and TeO_3 (and /or TeO_{3+1}) trigonal pyramids structure, respectively [24]. In our Li₂O modified in tungsten tellurite glasses, negligible shift of TeO₄ band is observed in Fig. 2. Moreover, the intensity of bands at 460 cm⁻¹ decrease which indicates the Te-O-Te linkage becoming increasing more depolymerized as TeO₄ is converted to TeO₃ with addition of Li₂O. Also in addition of Li₂O to tungsten tellurite (in between 10 and 15 mol%) results in a decrease of glass transition temperature due to the cleavage of the networks formed by TeO₄ trigonal bipyramid units and the increase of nonbridging oxygen (NBO) atoms [25], and W-O stretching shift also change.



Fig. 2 Raman spectrum of the undoped TWL glass

In the TWL glass, Er^{3+} ions are surrounded by NBO of TeO₄ tetrahedra and energy of local vibrational mode coupled to the Er^{3+} centers could be smaller than that of the maximum phonon energy of glass network. As per observed result the phonon energy of tungsten –tellurite glass is evidently larger than that of tellurite glass at 750 cm⁻¹. In our case change in energy transfer in Er^{3+} caused by the substitution of alkali ions can be neglected because maximum phonon energy are kept unchanged. The lifetime of the ⁴I_{11/2} level of Er^{3+} in TWL tellurite glasses may be determined almost entirely by the non-radiative decay rate [26]. Since energy gap of transition ⁴I_{11/2} \rightarrow ⁴I_{15/2} is about 3,600 cm⁻¹ then the non –radiative transition rate can be descried as

$$\tau^{-1} = W_{\text{non-rad}} \approx 10^7 \exp[-2(p-2.4)],$$
 (1)

where, $p=\Delta E/\hbar\omega$ is the number of phonon energy $\hbar\omega$ required to bridge the energy ΔE . Thus from the calculation of lifetime of the ${}^{4}I_{11/2}$ level of Er^{3+} in TWL glass is ~1.89 µs, and that of tellurite glass is 12 µs, and that of silica glass is close to 1 µs. In fact, the lifetime of ${}^{4}I_{11/2}$ level of Er^{3} ⁺ in TWL glass is comparable to that of silica glass, so it is possible to pump efficiently at 980 nm and excite ground level ${}^{4}I_{15/2}$ to upper level ${}^{4}I_{11/2}$ of Er^{3+} ions. Thus the enhanced green and the red radative transitions were observed from the reduced non-radiative transition in the TWL glass system. Further this TWL glass for 1.5 µm emission may be used by introduction of high phonon energy oxide such as P₂O₅, which increased the non-radiative decay rate of Er^{3+} ions and reduced the upconversion luminescence.

Upconversion Fluorescence and Mechanism

Figure 3 presents the upconversion emission spectra of the TWL: Er³⁺/Yb³⁺(2 mol%:2 mol%) glass sample in the wavelength range of 500-700 nm under 980 nm diode laser excitation at room temperature. The observed upconversion luminescence of intense green and red spectral bands has peaks at wavelength 540 nm with small shoulder at 525 nm, and 670 nm with shoulder 640 nm are simultaneously observed. The wavelength 540 nm and 525 nm corresponds to Er^{3+} : ${}^{4}\mathrm{S}_{3/2} \rightarrow {}^{4}\mathrm{I}_{15/2}$, ${}^{2}\mathrm{H}_{11/2} \rightarrow {}^{4}\mathrm{I}_{15/2}$ transition respectively while both 670 nm and 640 nm corresponding to ${}^{4}F_{9/2} \rightarrow {}^{4}I_{15/2}$ transition. With increasing the excitation intensity on measurement of emission intensity upconversion red emission intensity relatively increases to the green emission. However green emission is bright enough to be perceived with the naked eye even at low pump power (86 mW). It is also noted that red and green emission bands present considerably large half widths (5 nm around 520 nm, 16 nm around 560 nm, and 23 nm around 670 nm) due to the high

Fluorescence Intensity(A.U.)

100000

80000

60000

40000

20000

0

450

990mW

930mW

800mW

650mW

500mW

350mW 96mW

500

. 11/2

Fig. 3 Upconversion fluorescence spectra of TWL: Er (2 mol%)-Yb (2 mol%) glass at room temperature

550

600

Wavelength(nm)

650

700

⁴ F_{9/2}

local heating of the sample, caused by efficient 980 nm absorption by Yb^{3+} ions. These visible emissions not only are lower Stark components of the emitting levels ${}^{2}H_{11/2}$, ${}^{4}S_{3/2}$ and ${}^{4}F_{9/2}$ responsible for emission, but also upper ones, which are thermally populated.

To obtain the optimum doping concentration of Er³⁺ ions, a series of glasses doped with different Er³⁺ concentration were prepared, in which the Yb₂O₃ content were all 2 mol% fraction. When increasing the Er^{3+} concentration, the fluorescence intensity increases at the beginning due to an increasing absorption of pump energy and energy transfer from Yb^{3+} to Er^{3+} as shown in Fig. 4. On comparison the emission intensity for concentration of Er₂O₃ at 200mw pump power, Initially green emission intensity increase unto 1 mol% and then gradually red emission at 670 nm increase unto 2 mol%. The plot clearly shows that the glass containing 2 mol% of Er^{3+} ions is best upconversion emission intensity of TWL glass. And when Er^{3+} concentration is beyond a 2 mol% concentration, the intensity remain constant and/or decreases. This stronger red emission observed due to shortened distance between Er^{3+} ions at higher concentration. We assume that too much Yb³⁺ ions will cause the increase in back energy transfer from Er^{3+} to Yb^{3+} , i.e. ${}^{4}I_{11/2}(Er^{3+})+{}^{2}F_{7/2}(Yb^{3+}) \rightarrow {}^{4}F_{15/2}$ $(Er^{3+})+{}^{2}F_{5/2}(Yb^{3+})$. In addition, the up conversion intensities of Er³⁺ ions will become stronger with increasing Yb³⁺ ions in the TWL glass.

Figure 5 depicts the log-log dependence of the integrated green (540 nm) and red(670nm) emission intensities on the



Fig. 4 Upconversion luminescence intensities as a function of $\rm Er_2O_3 concentration$ at fix $\rm Yb_2O_3(2\ mol\%)$

excitation power at 980 nm. The Upconversion intensity I_{UP} will be proportional to m^{th} power of the I_{IR} excitation Intensity i.e. $I_{UP} \propto I_{IR}^m$, where, m is the number of I_{IR} photons required to convert a visible photon. A plot of ln (I_{UP}) versus ln (I_{IR}) yields a straight line with slope m of which the value is obtained from the slope of the graph [27]. Such dependence of integrated green and red emission intensities I_{UP} on excitation power intensity I_{RI} present in Fig. 5, which indicate that for the green and red emission for $\mathrm{Er}^{3+}/\mathrm{Yb}^{3+}$ (2 mol%:2 mol%) codoped lithium tungsten tellurite glass has given the straight lines with slop of 2 and 2.1 respectively. The results confirm that a two- photon upconversion process is assigned to the green and red



Fig. 5 Dependence of red and green upconversion fluorescence intensity on 980 nm LD excitation pump power for TWL glass

emission from, ${}^{4}S_{3/2}$ and ${}^{4}F_{9/2}$ level of Er^{3+} respectively. Based on the energy matching conditions and the dependence on excitation power, the possible upconversion mechanisms for the emission bands are discussed based on the energy level of Er³⁺ and Yb³⁺ ions presented in Fig. 6. Usually, the up-conversion emission of Er^{3+} can be explained by several well- known mechanisms such as excited -state absorption (ESA), energy transfer up-conversion (ETU) [27]. In singly Er^{3+} ions doped glass samples with 980 nm photons are responsible for excitation of the ${}^{4}I_{11/2}$ level which can also populate the lower-lying ${}^{4}I_{13/2}$ through multiphonon decay. Besides possibly originating near-infrared emission around 1550 nm and 2800 nm, both these metastable levels can undergo ESA of 980 nm photons leading to excitation of upper ${}^{4}F_{9/2}$ and ${}^{4}F_{7/2}$ levels. Given the small energy gap separating the levels ${}^4F_{7/2}$ and $({}^2H_{11/2}, \; {}^4S_{3/2}),$ the latter are rapidly populated via multiphonon decay from ${}^{4}F_{7/2}$. In ETU process involves not only two closely neighboring Er³⁺ ions but also in Er³⁺ and Yb³⁺ ions. In Er³⁺/Yb³⁺ co-doped samples the Yb \rightarrow Er energy transfers (ET₁, ET₂ and ET₃ in Fig. 6) are known to be the most significant contribution for 980 nm upconversion excitation. For the green emissions, in first step, the ⁴I_{11/2} level is directly excited with 980 nm light ground state absorption process (GSA) and /or by energy transfer (ET) process from ${}^{2}F_{5/2}$ level of Yb³⁺: ${}^{2}F_{5/2}(Yb^{3+})+$ ${}^{4}I_{15/2}(\text{Er}^{3+}) \rightarrow {}^{2}F_{7/2}(\text{Yb}^{3+}) + {}^{4}I_{11/2}(\text{Er}^{3+})$. Since Yb^{3+} has a much large absorption cross-section than Er³⁺ in the 980 nm region, the ET process is dominated to the excitation of ${}^{4}I_{11/2}$ level.

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Thus the second step involved the excitation process based on long-lived ${}^{4}I_{11/2}$ level as follows: cross-relaxation(CR), ${}^{4}I_{11/2}(Er^{3+}) + {}^{4}I_{11/2}(Er^{3+}) \rightarrow {}^{4}F_{7/2}(Er^{3+}) + {}^{4}I_{15/2}(Er^{3+});$ excited state absorption (ESA), ${}^{4}I_{11/2}(\text{Er}^{3+})+$ a photon $\rightarrow {}^{4}F_{7/2}(\text{Er}^{3+});$ ET, ${}^{2}F_{5/2}(Yb^{3+})+{}^{4}I_{11/2}(Er^{3+}) \rightarrow {}^{2}F_{7/2}(Yb^{3+})+{}^{4}F_{7/2}(Er^{3+})$. The populated ${}^4F_{7/2}$ level of \mbox{Er}^{3+} then relaxes rapidly and nonradiatively to the next lower level ${}^{2}H_{11/2}$ and ${}^{4}S_{3/2}$ resulting form the small energy gap between them. The ${}^{2}\text{H}_{11/2}$ state can also decay to the ${}^{4}\text{S}_{3/2}$ state due to multiphonon relaxation process. It is found that estimated energy gap between the ${}^{2}H_{11/2}$ and next lower level state ${}^{4}S_{3/2}$ is $\sim 800 \text{ cm}^{-1}$ [28], thus the multiphonon relaxation rate is very large and ${}^{2}H_{11/2} \rightarrow {}^{4}I_{15/2}$ transition is reduced. The above process produced the ${}^{4}S_{3/2} \rightarrow {}^{4}I_{15/2}$ transition for green emission centered at 540 nm. In our experiment with increasing Er_2O_3 concentration, the energy transfer rate increases, via Yb³⁺ and thus the visible emissions intensity increases. Therefore, the nonradiative transition rate through multiphonon relaxation from ${}^{4}S_{3/2}$ becomes smaller than from ${}^{2}\text{H}_{11/2}$, and the upconversion intensity from ${}^{4}\text{S}_{3/2} \rightarrow {}^{4}\text{I}_{15/2}$ transition is larger than from ${}^{2}H_{11/2} \rightarrow {}^{4}I_{15/2}$ transition as this result obtained for composition of 68TeO₂-15Li₂CO₃-13WO₃: 2Yb₂O₃-2Er₂O₃ glass.

The red upconversion emission is originated form the ${}^{4}F_{9/2} \rightarrow {}^{4}I_{15/2}$ transition centered at 670 nm and there exist two possible pumping mechanisms. The first pumping mechanism comprises the population of the ${}^{4}S_{3/2}$ state, followed by a fast nonradiative decay through multiphonon interaction from the population ${}^{4}S_{3/2}$ level to ${}^{4}F_{9/2}$ level and

Fig. 6 Partial energy level diagrams of Er^{3+} and Yb^{3+} . The pumping transition at 980 nm along with the up conversion emission in the green and red spectral region



then to the ground state. In other possible mechanism the population of ${}^{4}F_{9/2}$ is based on the processes as follows: ET from Yb³⁺, ${}^{2}F_{5/2}(Yb^{3+})+{}^{4}I_{13/2}(Er^{3+}) \rightarrow {}^{2}F_{7/2}(Yb^{3+}) +{}^{4}F_{9/2}$ (Er³⁺); cross relaxation between Er³⁺ ions, ${}^{4}I_{13/2} + {}^{4}I_{11/2} \rightarrow {}^{4}I_{15/2}$ $+{}^{4}F_{9/2}$; ESA, ${}^{4}I_{13/2}$ + a Photon $\rightarrow {}^{4}F_{9/2}$. The ${}^{4}I_{13/2}$ level is populated owing to the nonradiative relaxation from the upper ${}^{4}I_{11/2}$ level. Additional non-radiative decay form ${}^{4}S_{3/2}$ can also populate level ${}^{4}F_{9/2}$ favoring the red emission (${}^{4}F_{9/2} \rightarrow {}^{4}I_{15/2}$). At Er₂O₃ concentration is high enough that allow the interaction of these ions and various cross- relaxation energy transfers involving levels ${}^{4}I_{11/2}$ and ${}^{4}I_{13/2}$, playing a role in the excitation to the upper levels. Further relative increase of the red emission intensity to the green emission at higher 2 mol% Er^{3+} indicated the short distance between Er^{3+} ions benefit the red emission [29, 30]. In our TWL glass composition with higher concentration of Li₂O (15 mol%) may promote the formation of new impurity phase which might benefit the aggregation of Er^{3+} ions and the rate of cross relaxation energy transfer between Er³⁺ions would be improved resulting in the red emission being enhanced again. Thus the TWL is more preferable for red emission than green emission on pump power variation. However almost another oxide based Er^{3+} Yb³⁺ codoped tellurite glass matrix dominated green emission intensity upon 980 nm excitation [12]. This characteristics of Li₂O modified in tungsten tellurite glass has never been reported in other Er^{3+/}Yb³⁺ codoped tellurite based glass materials pumped at around 980 nm.

Conclusions

Er³⁺/Yb³⁺ co-doped modified tungsten tellurite by Li₂O glasses have been fabricated and characterized for absorption, Raman and upconvesion luminescence spectrum. The analysis of Raman spectrum confirms Li₂O enters as an intermediate between a network formerTeO₂ and a network modifier WO3, as a result W-O stretching-band shift has been observed in undoped TWL. The maximum phonon energy of TWL glass measured about 930 cm⁻¹, which leads to estimated lifetime of ${}^{4}I_{11/2}$ of Er³⁺ions~1.89 µs. With increasing Er₂O₃ at fix Yb₂O₃(2 mol%) in TWL glass matrix, the intensity green emission at 525 nm increase slightly and 560nm increases largely while red emission at 670 nm increase more than that of green at 2 mol% of Er₂O₃. Under 980 nm pump power variation for TWL :2Er-2Yb(mol%) glass sample shows relatively increase of intense red at 670 nm to green at 560 nm emissions indicate the quadratic dependent of upconversion intensity on pump intensity. The enhancement of red emission with increasing Er³⁺ ion concentration is due to cross-relaxation process. Thus it is found that TWL glass system may is considered for upconversion device although it is most prospective material for fiber and amplifier. The Judd-Ofelt intensity parameter and radiative lifetime estimation will be in progress in next work.

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